

Motion of One Electron in a Radiated $B^{(3)}$ Field
and a Static Magnetic Field.

1) Motion in a Static Magnetic Field

Angular frequency:

$$\Omega = \frac{eB}{m} \quad - (1)$$

Angular momentum:

$$\underline{J} = e r A \quad - (2)$$

Kinetic energy:

$$T = e^2 A^2 / (2m) \quad - (3)$$

2) Motion in the $B^{(3)}$ Radiated Field.

a) Direct Integration of Dirac Equation

Angular frequency:

$$\Omega = \frac{(\gamma m v + e A^{(0)}) c^2}{\left(r_0 + \frac{e A^{(0)}}{\gamma m c} \right) \left(m c^2 (1 + \gamma) + e \phi \right)} \quad - (4)$$

Angular Momentum:

- (5)

$$\underline{J} = \gamma \underline{r}_0 \times \underline{p} + e \underline{r}_0 \times \underline{A} + e \int \underline{A} dt \times \underline{v} + \frac{e^2}{m \gamma} \int \underline{A} dt \times \underline{A}$$

$$\underline{J} = \left(r_0 + \frac{e A^{(0)}}{\gamma m c} \right) (\gamma m v + e A^{(0)}) \quad - (5a)$$

2) Kinetic Energy:

$$T = \frac{(\gamma m v + e A^{(0)})^2 c^2}{m c^2 (1 + \gamma) + e \phi} \quad - (6)$$

b) Hamilton - Jacobi Method

Angular frequency:

$$\Omega = \left(\omega^2 + \frac{e^2 B^{(0)2}}{m^2} \right)^{1/2} \quad - (7)$$

Angular Momentum:

$$\underline{J}^{(3)} = \frac{e^2 c^2}{\omega^2} \left(\frac{B^{(0)}}{(m^2 \omega^2 + e^2 B^{(0)2})^{1/2}} \right) \underline{B}^{(3)} \quad (3)$$

Kinetic Energy:

$$\quad \quad \quad - (8)$$

$$T = \frac{1}{2} \Omega \underline{J}$$

$$= \frac{e^2 c^2}{2 m \omega^2} B^{(0)2} \quad - (9)$$

Notes

Ω = angular frequency of the electron

ω = angular frequency of the electromagnetic field

$$\gamma = \left(1 - \frac{u^2}{c^2} \right)^{-1/2}$$

v = velocity of electron, r = position of electron

$$\underline{r} = \underline{r}_0 + \frac{e}{m \gamma} \int \underline{A} dt$$

3) Notes on Static Magnetic Field

The angular momentum is:

$$\underline{J} = \underline{r} \times \underline{p} = e \underline{r} \times \underline{A} \quad - (10)$$

from the minimal prescription. In the standard model:

$$\underline{B} = \nabla \times \underline{A}, \quad \underline{A} = \frac{1}{2} \underline{B} \times \underline{r} \quad - (11)$$

so: $\underline{J} = e r A, \quad A = B r. \quad - (12)$

The kinetic energy is:

$$T = \frac{e^2 A^2}{2m} \quad - (13)$$

so the angular velocity is:

$$\Omega = \frac{dT}{dJ} = \frac{eA}{rm} = \frac{eB}{m}. \quad - (14)$$

The $B^{(3)}$ Field from Equation (5)

This is given by:

$$\underline{J}^{(3)} = \frac{e^2}{2\gamma} \int \underline{A}^{(1)} dt \times \underline{A}^{(2)} \quad - (15)$$

where: $\underline{A}^{(1)} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} - i\underline{j}) e^{i(\omega t - krz)} \quad - (16)$

$$\underline{A}^{(2)} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i(\omega t - krz)} \quad - (17)$$

so: $\underline{J}^{(3)} = \frac{e^2 A^{(0)2}}{m\gamma\omega} \underline{e}^{(3)} \quad - (18)$

We now we:

$$4) \quad A^{(0)} = \frac{c}{\omega} B^{(0)} \quad - (19)$$

and: $\underline{B}^{(3)} = B^{(0)} \underline{e}^{(3)} \quad - (20)$

to find: $\underline{J}^{(3)} = \frac{e^2 c^2}{\gamma m \omega^3} B^{(0)} \underline{B}^{(3)} \quad - (21)$

In the non-relativistic limit:

$$\gamma \rightarrow 1, \quad m\omega \gg eB^{(0)} \quad - (22)$$

and eqn. (21) becomes eqn (8) from Φ HJ method

For a static magnetic field:

$$\underline{J} = e r^2 B = e \underline{\Phi} \quad - (22)$$

where $\underline{\Phi} = r^2 B = \text{magnetic flux.}$

Therefore the quantum of flux is:

$$\underline{\Phi} = \frac{\hbar}{e} \quad (\text{i.e. } \underline{J} \rightarrow \hbar) \quad - (23)$$

From eq. (21):

$$\underline{J} = \frac{e^2 c^2}{\gamma m \omega^3} B^{(0)} = \frac{\mu_0 e^2 c}{\gamma m \omega} \left(\frac{\underline{I}}{\omega^2} \right)$$

5) Here:
$$I = \frac{c}{\mu_0} B^{(0)2} \quad - (24)$$

is the electromagnetic field's power density
(watts per square metre).

From a comparison of eqns. (22) and (23) it is seen that $\underline{B}^{(3)}$ and \underline{B} are completely different.

Induced Magnetic Dipole Moment

1) Static Magnetic Field

This is calculated from the gyromagnetic ratio

$$\underline{m} = \frac{e}{2m} \underline{J} \quad - (25)$$

$$= \left(\frac{e^2 r^2}{2m} \right) \underline{B} = \frac{e^2}{2m} \Phi \underline{k} \quad - (26)$$

So:
$$\chi = \frac{e^2 r^2}{2m} \quad - (27)$$

is the susceptibility.

2) The $\underline{B}^{(3)}$ Field

$$\underline{m}^{(3)} = \left(\frac{e^3 c^2}{28m^2 \omega^3} \right) B^{(0)} \underline{B}^{(3)} \quad - (28)$$

6)

Here:

$$\beta = \frac{e^3 c^2}{2\gamma m^2 \omega^3} \quad - (29)$$

is the magnetic hypermagnetizability.

Conclusions

- 1) A static magnetic field interacts with an electron through the static magnetic susceptibility (27) to induce a magnetic dipole moment.
- 2) The $\underline{B}^{(3)}$ field interacts with an electron through the frequency dependent hypermagnetizability (29) to induce a magnetic dipole moment.