

4.2(3): Calculation of the Newtonian Velocity in the Rotating Frame.

In the usual Newtonian theory:

$$v_N^2 = \left(\frac{dr}{dt}\right)^2 + r^2 \left(\frac{d\phi}{dt}\right)^2 \quad - (1)$$

in plane polar coordinates r and ϕ . For the conic section:

$$r = \frac{d}{1 + \epsilon \cos \phi} \quad - (2)$$

where d is the half right semi-major axis and ϵ the eccentricity.
From Lagrangian dynamics the constant angular momentum is:

$$L = m r^2 \frac{d\phi}{dt} \quad - (3)$$

Now we:

$$\frac{dr}{dt} = \frac{dr}{d\phi} \frac{d\phi}{dt} \quad - (4)$$

and

$$\frac{dr}{d\phi} = \frac{\epsilon r^2}{d} \sin \phi \quad - (5)$$

to find that

$$v_N^2 = \frac{L^2}{m} \left(\frac{1}{r^2} + \left(\frac{\epsilon}{d} \right) \sin^2 \phi \right) \quad - (6)$$

where

$$\sin^2 \phi = 1 - \cos^2 \phi \quad - (7)$$

and

$$\cos \phi = \frac{1}{\epsilon} \left(\frac{d}{r} - 1 \right) \quad - (8)$$

so

$$v_N^2 = \frac{L^2}{m} \left(\frac{1}{r^2} + \left(\frac{\epsilon}{d} \right)^2 \left(1 - \frac{1}{\epsilon^2} \left(\frac{d}{r} - 1 \right)^2 \right) \right)$$

$$= \frac{L^2}{m} \left(\frac{\epsilon^2 - 1}{d^2} + \frac{2}{dr} \right) \quad - (9)$$

Finally, we:

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Finally, we:

$$2) \quad L^2 = m^2 M G d \quad - (10)$$

to find that:

$$v_n^2 = M G \left(\frac{2}{r} + \frac{e^2 - 1}{d} \right)$$

$$= M G \left(\frac{2}{r} - \frac{1}{a} \right) \quad - (11)$$

where

$$\frac{1}{a} = \frac{1 - e^2}{d} \quad - (12)$$

is the semi major axis.

In the rotating frame:

$$v_n^2 = M G \left(\frac{2}{r} - \frac{1}{a'} \right) \quad - (13)$$

where

$$a' = \frac{d'}{1 - e'^2} \quad - (14)$$

Here:

$$d' = \frac{L'^2}{m^2 M G} \quad - (15)$$

where:

$$L' = \mu r^2 \omega'$$

$$= \mu r^2 \frac{d\phi'}{dt} = \mu r^2 \frac{d}{dt} (\phi + \omega_1 t)$$

$$= \mu r^2 \left(\omega + \omega_1 + t \frac{d\omega_1}{dt} \right) \quad - (16)$$

= constant of motion.

Also,

$$e'^2 = 1 + \frac{2H' L'^2}{\mu (m M G)^2} \quad - (17)$$

where the Hamiltonian H' in the rotating frame is:

$$H' = \frac{1}{2} \mu \left(\frac{dr}{dt} \right)^2 + \frac{1}{2} \frac{L'^2}{mr^2} - \frac{nM\Gamma}{r} \quad (18)$$

The Hamiltonian is constant of motion, as follows:

$$\begin{aligned} H' &= \frac{1}{2} m v'^2 - \frac{nM\Gamma}{r} \\ &= \frac{1}{2} m M \Gamma \left(\frac{2}{r} - \frac{1}{a'} \right) \quad (19) \\ &= -\frac{nM\Gamma}{2a'} \end{aligned}$$

so a' is also a constant of motion.

In order to compute the Hamiltonian use:

$$r = \frac{d'}{1 + \epsilon' \cos \phi'} \quad (20)$$

and

$$\begin{aligned} \frac{dr}{dt} &= \frac{-\epsilon' d' \frac{d}{dt} \cos \phi'}{(1 + \epsilon' \cos \phi')^2} \quad (21) \\ &= \frac{\epsilon' d' \sin \phi'}{(1 + \epsilon' \cos \phi')^2} \frac{d\phi'}{dt} \end{aligned}$$

in eq. (18). Here:

$$\phi' = \phi + \omega_1 t, \quad (22)$$

and

$$\frac{d\phi'}{dt} = \frac{d\phi}{dt} + \omega_1 + \frac{d\omega_1}{dt} \quad (23)$$

+) These equations allow H' and L' to be calculated in terms of a and ω . For an orbit:

$$t = T' - (24)$$

where:

$$T'^2 = \frac{4\pi^2 \mu}{mM\dot{G}} a'^3 - (25)$$

in the rotating frame. Eq (25) is Kepler's third law. Having calculated the Hamiltonian H' , the semi major axis in the rotating frame is calculated from eq. (19). This is a' . Knowing a' , the time taken for one orbit in the rotating frame, T' , can be calculated from Kepler's third law in the rotating frame, eq. (25).

In the rotating frame, d' and e' are the observed d' and e' . For the Hulse Taylor binary pulsar and the S2 star, d' and e' are known and have been used in previous UFT papers. So v_N' can be calculated in the rotating frame and compared with v_N in the usual Newtonian theory. This gives the effect of rotation on the Newtonian linear velocity, i.e. the orbital linear velocity.

Similarly, the effect of frame rotation can be found for d , e , T , L and H , and also for the Lagrangian dynamics is changed by ω .

5) Note carefully that a universal law of precession:

$$\Delta\phi = \frac{2\pi}{2} \left(v_N^2 + r^2 (\omega_1^2 + 2\omega\omega_1) \right) - (26)$$

the Newtonian velocity v_N is the $\frac{d\mathbf{r}}{dt}$ of the

unrotated frame, i.e. by:

$$v_N^2 = \dot{r}^2 \left(\frac{2}{r} - \frac{1}{a} \right) - (27)$$

with:

$$a = \frac{\alpha}{1-\epsilon^2} - (28)$$

and

$$\alpha = \frac{L^2}{m^2 MG} - (29)$$

$$\epsilon^2 = 1 + \frac{2HL^2}{\mu(mMG)^2} - (30)$$

$$H = \frac{1}{2} \mu \left(\frac{dr}{dt} \right)^2 + \frac{1}{2} \frac{L^2}{mr^2} - \frac{mMG}{r} - (31)$$

and

$$\frac{dr}{dt} = \frac{\epsilon \sin\phi}{(1+\epsilon\cos\phi)^2} \frac{d\phi}{dt} - (32)$$

with

$$\omega = \frac{d\phi}{dt}$$
