

312(5): The Density of States due to Photon Mass

In previous work the density of states was defined by:

$$\frac{dN}{V} = \frac{1}{3c^3 \pi^2} \left((\omega + d\omega)^3 - \omega_0^3 \right)^{3/2} \quad - (1)$$

In these expressions:

$$(\omega^2 - \omega_0^2)^{3/2} = \left(\omega^2 \left(1 - \frac{\omega_0^2}{\omega^2} \right) \right)^{3/2}$$
$$= \omega^3 \left(1 - \frac{\omega_0^2}{\omega^2} \right)^{3/2} \quad - (2)$$

If $\omega_0 \ll \omega$ - (3)

then $\left(1 - \frac{\omega_0^2}{\omega^2} \right)^{3/2} \sim 1 - \frac{3}{2} \left(\frac{\omega_0}{\omega} \right)^2 \quad - (4)$

Similarly:

$$((\omega + d\omega)^2 - \omega_0^2)^{3/2} = (\omega + d\omega)^3 \left(1 - \frac{\omega_0^2}{(\omega + d\omega)^2} \right)^{3/2}$$
$$\sim (\omega + d\omega)^3 \left(1 - \frac{3}{2} \frac{\omega_0^2}{(\omega + d\omega)^2} \right)$$
$$= (\omega + d\omega)^3 - \frac{3}{2} \omega_0^2 (\omega + d\omega) \quad - (5)$$

Therefore:

$$\frac{dN}{V} \sim \frac{1}{3c^3 \pi^2} \left((\omega + d\omega)^3 - \frac{3}{2} \omega_0^2 (\omega + d\omega) - \omega^3 \left(1 - \frac{3}{2} \left(\frac{\omega_0}{\omega} \right)^2 \right) \right) \quad - (6)$$

$$= \frac{1}{3c^3\pi^2} \left((\omega + d\omega)^3 - \omega^3 - \frac{3}{2} \omega_0^2 (\omega + d\omega) + \frac{3}{2} \omega^3 \left(\frac{\omega_0}{\omega} \right)^2 \right) \quad (7)$$

$$= \frac{1}{3c^3\pi^2} \left((\omega + d\omega)^3 - \omega^3 + \frac{3}{2} (\omega \omega_0^2 - \omega_0^2 \omega - \omega_0^2 d\omega) \right)$$

Therefore:

$$\frac{dN}{V} = \frac{1}{3c^3\pi^2} \left((\omega + d\omega)^3 - \omega^3 \right) - \frac{\omega_0^2}{2c^3\pi^2} d\omega \quad (8)$$

The first term on the RHS of eq. (8) is the usual Rayleigh-Jeans density of states but the photon rest mass m_0 contributes the term:

$$\boxed{\frac{dN}{V} = - \frac{\omega_0^2}{2c^3\pi^2} d\omega} \quad (9)$$

where the photon rest frequency is:

$$\omega_0 = m_0 c^2 / \hbar \quad (10)$$

the energy density due to eq. (9) is:

$$\frac{dE}{V} = - \frac{\omega_0^2}{2c^3\pi^2} \langle \hbar\omega \rangle d\omega \quad (11)$$

where

$$\langle \hbar\omega \rangle = \frac{\hbar\omega}{e^y - 1} \quad (12)$$

with

$$y = \frac{\hbar\omega}{kT} \quad (13)$$

3) In the high temperature approximation:
 $\hbar\omega \ll kT$ — (14)

then
$$\frac{dE}{V} = - \frac{\omega_0^2 kT}{2c^3 \pi^2} d\omega$$
 — (15)

Therefore the correction to the Rayleigh Jeans law due to photon rest mass leads to the energy density correction:

$$\frac{\bar{E}}{V} = - \left(\frac{\omega_0^2 kT}{2c^3 \pi^2} \right) \omega$$
 — (16)

by integrating eq. (15) Therefore:

$$\boxed{\frac{\bar{E}}{V} = - \left(\frac{\omega_0^2 kT}{2c^3 \pi^2} \right) \omega}$$
 — (17)
 for $\hbar\omega \ll kT$

For black body radiation:

$$\bar{E} = - \left(\frac{\omega_0^2 kT}{2c^3 \pi^2} \right) \int_0^\infty d\omega$$
 — (18)

→ ∞
 Therefore the photon rest mass affects the Stefan Boltzmann law very considerably if high enough frequencies are taken into account.

4) The rigorous calculation of the cavities is:

$$\bar{\Phi} = \frac{cE}{V} = \frac{-\omega_0^2}{2c^2\pi^2} \int \frac{\hbar\omega}{e^{\gamma}-1} d\omega \quad - (19)$$

and for black body radiation:

$$\bar{\Phi} = - \frac{\hbar\omega_0^2}{2c^2\pi^2} \int_0^\infty \frac{\omega}{e^{\gamma}-1} d\omega \quad - (20)$$

The integral can be worked out with:

$$\int_0^\infty \left(\frac{x^{2n-1}}{e^x-1} \right) dx = (2\pi)^{2n} \frac{B_n}{4n} \quad - (21)$$

where $B_1 = \frac{1}{6}$, $B_2 = \frac{1}{30}$, $B_3 = \frac{1}{42}$, ...

are Bessel functions

The important thing to note is that the cavities due to photon rest mass of a Stefan Boltzmann law becomes very large when:

$$\hbar\omega \ll kT \quad - (22)$$

and at very high frequencies ω . Otherwise for eq. (20) it is very small.

For example, if two gamma rays frequencies ω_1 and ω_2 define a range of ray frequencies then for eq. (7):

$$\underline{\Phi} = - \left(\frac{\omega_0^2 k T}{2 c^3 \pi^2} \right) (\omega_1 - \omega_2) \quad - (23)$$

ii which:

$$k = 1.38066 \times 10^{-23} \text{ J K}^{-1} \quad - (24)$$

and

$$c = 2.997925 \times 10^8 \text{ m s}^{-1} \quad - (25)$$