

## 254(3) : Vectorial Format for Hodge Duals

Consider the Cartan identity:

$$D_\mu T_{\nu\rho}^a + D_\rho T_{\mu\nu}^a + D_\nu T_{\rho\mu}^a := R_{\mu\nu\rho}^a + R_{\rho\mu\nu}^a + R_{\nu\rho\mu}^a - (1)$$

The Hodge dual identity is:  $\downarrow$

$$D_\mu \tilde{T}^{a\mu\nu} := \tilde{R}_\mu^{a\mu\nu} - (2)$$

Therefore the space part of eq. (1) is Hodge transformed as follows:  $- (3)$

$$D_1 T_{23}^a + D_3 T_{12}^a + D_2 T_{31}^a := R_{123}^a + R_{312}^a + R_{231}^a$$

$\downarrow$

$$D_1 \tilde{T}^{a10} + D_2 \tilde{T}^{a20} + D_3 \tilde{T}^{a30} := \tilde{R}_1^{a10} + \tilde{R}_2^{a20} + \tilde{R}_3^{a30} - (4)$$

In vector notation:

$$\underline{\nabla} \cdot \underline{T}^a + \underline{\omega}^a_b \cdot \underline{T}^b := \underline{v}^b \cdot \underline{R}^a_b - (5)$$

$\downarrow$

$$\underline{\nabla} \cdot \underline{\tilde{T}}^a + \underline{\omega}^a_b \cdot \underline{\tilde{T}}^b = \underline{v}^b \cdot \underline{\tilde{R}}^a_b - (6)$$

The reverse process is:

$$D_\mu \tilde{T}_{\nu\rho}^a + D_\rho \tilde{T}_{\mu\nu}^a + D_\nu \tilde{T}_{\rho\mu}^a := \tilde{R}_{\mu\nu\rho}^a + \tilde{R}_{\rho\mu\nu}^a + \tilde{R}_{\nu\rho\mu}^a - (7)$$

$\downarrow$

$$D_\mu \tilde{T}^{a\mu\nu} := \tilde{R}_\mu^{a\mu\nu} - (8)$$

2) i.e.

$$\underline{\nabla} \cdot \underline{\tilde{T}}^a + \underline{\omega}^a_b \cdot \underline{\tilde{T}}^b = \underline{v}^b \cdot \underline{\tilde{R}}^a_b - (9)$$

$$\underline{\nabla} \cdot \underline{T}^a + \underline{\omega}^a_b \cdot \underline{T}^b + \underline{v}^b \cdot \underline{R}^a_b - (10)$$

For each index  $a$  define the electromagnetic field tensor as:

$$F^{\mu\nu} = \begin{pmatrix} 0 & -E^1 & -E^2 & -E^3 \\ E^1 & 0 & -cB^3 & cB^2 \\ E^2 & cB^3 & 0 & -cB^1 \\ E^3 & -cB^2 & cB^1 & 0 \end{pmatrix} - (11)$$

then eqs. (5) and (6) produce:

$$\underline{\nabla} \cdot \underline{B}^a = 0 - (12)$$

with condition:

$$\underline{v}^b \cdot \underline{R}^a_b = \underline{\omega}^a_b \cdot \underline{T}^b - (13)$$

and eqs. (9) and (10) produce:

$$\underline{\nabla} \cdot \underline{E}^a = \frac{\rho}{\epsilon_0} - (14)$$

As shown in notes 254(1) and 254(2), eq. (5) reduces to:

$$\underline{\nabla} \cdot \underline{v}^b \times \underline{\omega}^a_b = 0 - (15)$$

and eq. (13) also reduces to eq. (15). The magnetic field is defined by:

$$\underline{B}^a = \underline{\nabla} \times \underline{A}^a - \underline{\omega}^a{}_b \times \underline{A}^b - (16)$$

with:

$$\underline{A}^a = A^{(0)} \underline{v}^a - (17)$$

so eq. (15) means that:

$$\underline{B}^a = \underline{\nabla} \times \underline{A}^a - (18)$$

Therefore:

$$\underline{\nabla} \cdot \underline{B}^a = \underline{\nabla} \cdot \underline{\nabla} \times \underline{A}^a = 0 - (19)$$

which is eq. (12), Q.E.D.

It is important to note that eqs. (5)  
and (6) describe the magnetic field, and eqs.  
(9) and (10) describe the electric field. This result is by  
definition (11). The space part of (11) is considered  
 in eq. (5) and it is magnetic. The Hodge dual  
 of eq. (6) also refers to the magnetic field. In  
 eq. (9) the space part of the Hodge dual of eq. (11)  
 is considered. The Hodge dual of eq. (11)  
 at mixed indices is, for each  $a$ :

$$\tilde{F}^{\mu\nu} = \begin{bmatrix} 0 & -cB^1 & -cB^2 & -cB^3 \\ cB^1 & 0 & E^3 & -E^2 \\ cB^2 & -E^3 & 0 & E^1 \\ cB^3 & E^2 & -E^1 & 0 \end{bmatrix} - (20)$$

4) So eqs. (9) and (10) give again eq. (15). However the electric field is defined by:

$$\underline{E}^a = -\underline{\nabla} \phi^a - \frac{\partial \underline{A}^a}{\partial t} - \omega^a_{\phantom{a}0b} \underline{A}^b + \omega^a_{\phantom{a}0b} \phi^b \quad (21)$$

So eq. (15) does not imply that the divergence of  $\underline{E}^a$  is zero in general. So we write at eq. (14) in general.

The next note will continue this development for  $n = 1, 2, 3$  in eqs. (2) or (8).

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