

225(1): Development of Electroweak Theory from the Minimal Prescription and Fermion Equation.

In the notation of UFT 172 of the fermion equation is:

$$\gamma_\mu \psi \sigma^\mu = mc \sigma^1 \psi \quad (1)$$

and can be written as two simultaneous equations:

$$(\hat{E} + c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^L = mc^2 \phi^R \quad (2)$$

$$(\hat{E} - c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^R = mc^2 \phi^L \quad (3)$$

where  $\phi^R = \begin{bmatrix} \psi_1^R \\ \psi_2^R \end{bmatrix}$ ,  $\phi^L = \begin{bmatrix} \psi_1^L \\ \psi_2^L \end{bmatrix} \quad (4)$

From eq. (3):

$$\phi^R = \left( \frac{mc^2}{\hat{E} - c \underline{\sigma} \cdot \underline{\hat{p}}} \right) \phi^L \quad (5)$$

So

$$(\hat{E} + c \underline{\sigma} \cdot \underline{\hat{p}})(\hat{E} - c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^R = m^2 c^4 \phi^R \quad (6)$$

$$(\hat{E} - c \underline{\sigma} \cdot \underline{\hat{p}})(\hat{E} + c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^L = m^2 c^4 \phi^L \quad (7)$$

As in note 173(4) the fermion equation is:

$$2) \sigma^0 \hat{E} \psi \sigma^0 - c \sigma^3 (\hat{p}_x \psi \sigma^1 + \hat{p}_y \psi \sigma^2 + \hat{p}_z \psi \sigma^3) = mc^2 \sigma^1 \psi \quad - (8)$$

and

$$\psi = \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} \quad - (9)$$

$$- (10)$$

The Pauli matrices are:

$$\sigma^0 = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \sigma^1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \sigma^2 = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \sigma^3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}.$$

The basic operators of quantum mechanics are:

$$\hat{E} = i\hbar \frac{\partial}{\partial t}, \quad \hat{p} = -i\hbar \nabla \quad - (11)$$

Eqs. (2) and (3) may be written as:

$$(\hat{E}^2 - c^2 \underline{\sigma} \cdot \underline{p} \underline{\sigma} \cdot \underline{p}) \phi^R = mc^4 \phi^R \quad - (12)$$

$$(\hat{E}^2 - c^2 \underline{\sigma} \cdot \underline{p} \underline{\sigma} \cdot \underline{p}) \phi^L = mc^4 \phi^L \quad - (13)$$

In note 173(4) the semi operator method of solving these equations was introduced:

$$\hat{E}^2 = \hat{E} \hat{E} = i\hbar \hat{E} \frac{\partial}{\partial t} \quad - (14)$$

where

$$E = \gamma mc^2 \quad - (15)$$

hence eq. (12) becomes:

$$3) (\gamma mc^2 \hat{E} - c^2 \underline{\sigma} \cdot \underline{p} \underline{\sigma} \cdot \underline{p}) \phi^R = m^2 c^4 \phi^R - (16)$$

$$\therefore \hat{E} \phi^R = H \phi^R - (17)$$

$$\text{where } H = \frac{1}{E} (c^2 \underline{\sigma} \cdot \underline{p} \underline{\sigma} \cdot \underline{p} + m^2 c^4) - (18)$$

is the Hamiltonian, "which":

$$E = \gamma mc^2, \quad \underline{p} = \gamma m \underline{v}, - (19)$$

$$\text{so } H = \frac{\gamma}{m} \underline{\sigma} \cdot \underline{v} \underline{\sigma} \cdot \underline{v} + \frac{mc^2}{\gamma} - (20)$$

So the Dirac equation becomes:

$$i \hbar \frac{\partial \phi^R}{\partial t} = H \phi^R - (21)$$

$$i \hbar \frac{\partial \phi^L}{\partial t} = H \phi^L - (22)$$

with H given by eq. (20).

Eq.s. (21) and (22) are equations for the right and left handed free electron.

Interaction with the electromagnetic field

is developed with the minimal prescription:

$$p^\mu \rightarrow p^\mu - e A^\mu - (23)$$



4) where  $p^\mu = \left( \frac{E}{c}, \underline{p} \right) \quad - (24)$

so:  $E \rightarrow E - e\phi, \quad - (25)$

$\underline{p} \rightarrow \underline{p} - e \underline{A}. \quad - (26)$

In FCE theory, these equations are true for each  
case of polarization a.

Defining  $\underline{\pi} = \underline{p} - e \underline{A} \quad - (27)$

eq. (16) becomes:

$\left( (E - e\phi)(\hat{E} - e\phi) - c^2 \underline{\sigma} \cdot \underline{\pi} \underline{\sigma} \cdot \underline{\pi} \right) \phi^R = m^2 c^4 \phi^R \quad - (28)$

so  $\hat{E} \phi^R = H \phi^R \quad - (29)$

$- (30)$

where  $H = \frac{m^2 c^4}{E - e\phi} + e\phi + c^2 \frac{\underline{\sigma} \cdot \underline{\pi} \underline{\sigma} \cdot \underline{\pi}}{E - e\phi}$

$= \frac{m^2 c^4}{E - e\phi} + e\phi + H_3$

5) The correctly relativistic Lande' and Thomas factors are obtained from:

$$H_3 = c^2 \underline{\sigma} \cdot \underline{\pi} (E - e\phi)^{-1} \underline{\sigma} \cdot \underline{\pi} \quad - (31)$$

$$= \frac{c^2}{E} \underline{\sigma} \cdot \underline{\pi} \left(1 - \frac{e\phi}{E}\right)^{-1} \underline{\sigma} \cdot \underline{\pi}$$

At this point an approximation is made:

$$\left(1 - \frac{e\phi}{E}\right)^{-1} \sim 1 + \frac{e\phi}{E} \quad - (32)$$

so it is assumed that:

$$e\phi \ll E \quad - (33)$$

where

$$E = \gamma m c^2 \quad - (34)$$

So:

$$H_3 = \frac{1}{\gamma m} (\underline{\sigma} \cdot \underline{\pi})(\underline{\sigma} \cdot \underline{\pi}) \quad - (35)$$

$$+ \frac{e}{\gamma^2 m^2 c^2} \underline{\sigma} \cdot \underline{\pi} \phi \underline{\sigma} \cdot \underline{\pi}$$

The first term gives the relativistic Lande' factor and the second term the relativistic Thomas factor.



6) Eq. (28) can be written as:

$$\left( E - e\phi - \frac{c^2 \underline{\sigma} \cdot \underline{\pi} \underline{\sigma} \cdot \underline{\pi}}{\hat{E} - e\phi} \right) \phi^R = \left( \frac{m^2 c^4}{E - e\phi} \right) \phi^R \quad (36)$$

Add  $mc^2$  to each side:

$$\left( E + mc^2 - e\phi - \frac{c^2 \underline{\sigma} \cdot \underline{\pi} \underline{\sigma} \cdot \underline{\pi}}{\hat{E} - e\phi} \right) \phi^R = \left( \frac{m^2 c^4}{E - e\phi} + mc^2 \right) \phi^R \quad (37)$$

In the non-relativistic approximation:  
 $E \rightarrow mc^2 \quad (38)$

so eq. (37) becomes:

$$\begin{aligned} & (mc^2 (2mc^2 - e\phi)) \phi^R = \\ & ((2mc^2 - e\phi)(\hat{E} - e\phi) - c^2 (\underline{\sigma} \cdot \underline{\pi})(\underline{\sigma} \cdot \underline{\pi})) \phi^R \\ & = (mc^2 (2mc^2 - e\phi)) \phi^R \quad (39) \end{aligned}$$

$$\therefore \hat{E} \phi^R = H \phi^R \quad (40)$$

where:

$$H = mc^2 + e\phi + c^2 (\underline{\sigma} \cdot \underline{\pi})(2mc^2 - e\phi)^{-1} (\underline{\sigma} \cdot \underline{\pi}) \quad (41)$$

7) The Hamiltonian is approximated by:

$$H \sim mc^2 + e\phi + \frac{1}{2m} \underline{\sigma} \cdot \underline{\pi} \left( 1 - \frac{e\phi}{2mc^2} \right) \underline{\sigma} \cdot \underline{\pi}$$

$$= mc^2 + e\phi + \frac{1}{2m} \underline{\sigma} \cdot \underline{\pi} \underline{\sigma} \cdot \underline{\pi}$$

$$+ \frac{1}{4m^2 c^2} \underline{\sigma} \cdot \underline{\pi} \phi \underline{\sigma} \cdot \underline{\pi} \quad - (42)$$

where  $\underline{\pi} = \underline{p} - e\underline{A} \quad - (43)$

The Landé factor, Zeeman effect, ESR, NMR, MR  $\underline{\pi}$  and FSR are obtained from:

$$H_3 = \frac{1}{2m} \underline{\sigma} \cdot \underline{\pi} \underline{\sigma} \cdot \underline{\pi} \quad - (44)$$

and Thomas factor, spin-orbit coupling and Darwin term from:

$$H_4 = \frac{1}{4m^2 c^2} \underline{\sigma} \cdot \underline{\pi} \phi \underline{\sigma} \cdot \underline{\pi} \quad - (45)$$

In the next note these terms will be developed for electromagnetic and weak field interaction

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