

223(7) : Cross Check & Evaluation of Relativistic Force.

Consider the velocity and acceleration in cylindrical polar coordinates :

$$\underline{v} = \underline{\dot{r}} = \frac{dr}{dt} = \frac{d}{dt} (r \underline{e}_r) \quad - (1)$$

$$= \dot{r} \underline{e}_r + r \underline{\dot{e}}_r,$$

$$\underline{a} = \frac{d\underline{v}}{dt} = (\ddot{r} - r\dot{\theta}^2) \underline{e}_r + (r\ddot{\theta} + 2\dot{r}\dot{\theta}) \underline{e}_\theta. \quad - (2)$$

The relativistic angular momentum is :

$$L = \gamma m r^2 \dot{\theta} \quad - (3)$$

and the orbit is :

$$r = \frac{d}{1 + E \cos \theta} \quad - (4)$$

Therefore

$$\frac{dr}{d\theta} = \frac{E}{d} r^2 \sin \theta \quad - (5)$$

and

$$\frac{dr}{dt} = \frac{dr}{d\theta} \frac{d\theta}{dt} = \frac{EL}{\gamma m d} \sin \theta. \quad - (6)$$

So:

$$\dot{r} = \frac{1}{\gamma} \left(\frac{LE}{md} \right) \sin \theta, \quad \dot{\theta} = \frac{L}{\gamma m r^2} \quad - (7)$$

It follows that

$$\begin{aligned} \ddot{r} &= \frac{1}{\gamma} \left(\frac{LE}{md} \right) \frac{d}{dt} \sin \theta \\ &= \frac{1}{\gamma} \left(\frac{LE}{md} \right) \cos \theta \frac{d\theta}{dt} \end{aligned} \quad - (8)$$

2) So
$$\ddot{r} = \left(\frac{L}{\gamma m r} \right)^2 \frac{e \cos \theta}{d} \quad - (9)$$

Also:

$$\begin{aligned} \ddot{\theta} &= -\frac{2L}{\gamma m r^3} \frac{1}{\gamma} \left(\frac{L e}{m d} \right) \sin \theta \\ &= -2 \left(\frac{L}{\gamma m} \right)^2 \frac{e}{d} \frac{1}{r^3} \sin \theta \quad - (10) \end{aligned}$$

So:
$$\ddot{r} - r \dot{\theta}^2 = \left(\frac{L}{\gamma m r} \right)^2 \left(\frac{e}{d} \cos \theta - \frac{1}{r} \right) \quad - (11)$$

$$r \ddot{\theta} + 2 \dot{r} \dot{\theta} = 0 \quad - (12)$$

Hence:

$$\underline{a} = \left(\frac{L}{\gamma m r} \right)^2 \left(\frac{e}{d} \cos \theta - \frac{1}{r} \right) \underline{e}_r \quad - (13)$$

where
$$\cos \theta = \frac{1}{e} \left(\frac{d}{r} - 1 \right) \quad - (14)$$

It follows that:

$$\boxed{\underline{a} = -\frac{1}{d} \left(\frac{L}{\gamma m r} \right)^2 \underline{e}_r} \quad - (15)$$

This is the relativistic acceleration due to eq. (4).

3) The relativistic force is defined by :

$$\underline{F} = m \frac{d}{dt} (\gamma \underline{v}) \quad - (16)$$

So

$$F = \frac{d}{dt} (\gamma m v) = m \left(\frac{d\gamma}{dt} v + \gamma \frac{dv}{dt} \right) \quad - (17)$$

where :

$$\frac{d\gamma}{dt} = \frac{d\gamma}{dv} \frac{dv}{dt} \quad - (18)$$

$$= \left(\frac{d}{dv} \left(1 - \frac{v^2}{c^2} \right)^{-1/2} \right) \frac{dv}{dt}$$

$$= \frac{v}{c^2} \frac{dv}{dt} \gamma^3$$

So

$$F = m \left(\frac{v^2}{c^2} \frac{dv}{dt} \gamma^3 + \gamma \frac{dv}{dt} \right)$$

$$= m \gamma \frac{dv}{dt} \left(\gamma^2 \frac{v^2}{c^2} + 1 \right)$$

$$= m \gamma \frac{dv}{dt} \left(\frac{v^2}{c^2} \left(1 - \frac{v^2}{c^2} \right)^{-1} + 1 \right)$$

$$F = m \gamma^3 \frac{dv}{dt} = m \gamma^3 a \quad - (19)$$

4) From eqs. (15) and (1a):

$$\underline{F} = -\frac{\gamma}{d} \left(\frac{L}{mr} \right)^2 \underline{e}_r \quad - (20)$$

and this is the same as eq. (17) of note 223 (6).

Conclusion

The relativistic force needed to describe the ellipse (4) is

$$\boxed{F^* = \gamma F} \quad - (21)$$

where F is the non-relativistic force, i.e.

$$F^* = \left(1 - \frac{v^2}{c^2} \right)^{-1/2} F \quad - (22)$$

or

$$\boxed{\begin{aligned} F^* &= \left(1 - \frac{v^2}{c^2} \right)^{-3/2} m \frac{dv}{dt} \\ &= \left(1 - \frac{v^2}{c^2} \right)^{-3/2} m \frac{d^2 r}{dt^2} \end{aligned}}$$