

179(1): The Generally Covariant Fermi Equation.

This equation can be derived straight forwardly from the Dirac wave equation:

$$(\square + R)\phi = 0 \quad - (1)$$

where ϕ is a tetrad. The fermi equation is the limit:

$$R \rightarrow \left(\frac{mc}{\hbar}\right)^2 \quad - (2)$$

Conversely, the fermi equation can be generalized by:

$$mc^2 \rightarrow \left(\frac{\hbar^2}{m}\right) R, \quad - (3)$$

where R is positive valued. Therefore the generally covariant fermi equation is:

$$(\underline{E} - \underline{V} - c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^R = \frac{\hbar^2}{m} R \phi^L \quad - (4)$$

$$(\underline{E} - \underline{V} + c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^L = \frac{\hbar^2}{m} R \phi^R \quad - (5)$$

Defining the kinetic energy as:

$$\underline{E} = \underline{E} - \frac{\hbar^2}{m} R \quad - (6)$$

Here is the notation of UFT 178:

$$(\underline{E} - \underline{V}) \phi_s^R = c \underline{\sigma} \cdot \underline{\hat{p}} \phi_s^L \quad - (7)$$

$$\left(\underline{E} - \underline{V} + 2 \frac{\hbar^2 R}{m}\right) \phi_s^L = c \underline{\sigma} \cdot \underline{\hat{p}} \phi_s^R \quad - (8)$$

Eqs. (7) and (8) give:

$$2) \quad (\epsilon - V) \phi_s^R = \frac{mc^2}{2\hbar^2 R} \underline{\sigma} \cdot \underline{\hat{p}} \left(\frac{\underline{\sigma} \cdot \underline{\hat{p}} \phi_s^R}{1 + \frac{m(\epsilon - V)}{2\hbar^2 R}} \right) - (9)$$

$$\text{If: } \frac{m(\epsilon - V)}{2\hbar^2 R} \ll 1 \quad - (10)$$

Then:

$$(\epsilon - V) \phi_s^R = \frac{mc^2}{2\hbar^2 R} \underline{\sigma} \cdot \underline{\hat{p}} \left(\left(1 - \frac{m(\epsilon - V)}{2\hbar^2 R} \right) \underline{\sigma} \cdot \underline{\hat{p}} \phi_s^R \right) - (11)$$

In eq. (11), R is defined by the tetrad postulate:

$$D_\mu \gamma_\nu^a = \partial_\mu \gamma_\nu^a + \omega_{\mu b}^a \gamma_\nu^b - \Gamma_{\mu\nu}^\lambda \gamma_\lambda^a = 0. \quad - (12)$$

$$\text{Define: } \omega_{\mu\nu}^a = \omega_{\mu b}^a \gamma_\nu^b \quad - (13)$$

$$\Gamma_{\mu\nu}^a = \Gamma_{\mu\nu}^\lambda \gamma_\lambda^a \quad - (14)$$

$$\text{Then } D_\mu \gamma_\nu^a = \Gamma_{\mu\nu}^a - \omega_{\mu\nu}^a \quad - (15)$$

$$\text{and } \square \gamma_\nu^a = \partial^\mu (\Gamma_{\mu\nu}^a - \omega_{\mu\nu}^a) \quad - (16)$$

3) The scalar curvature R is defined by :

$$\square \psi^a = -R \psi^a = \partial^\mu (\Gamma_{\mu\nu}^a - \omega_{\mu\nu}^a) \quad (17)$$

i.e

$$R = \psi^a \partial^\mu (\omega_{\mu\nu}^a - \Gamma_{\mu\nu}^a) \quad (18)$$

In general, R is a function of distance and of time, and changes under the influence of gravitation.

In the absence of gravitation :

$$R \rightarrow \left(\frac{mc}{\hbar} \right)^2 \quad (19)$$

In generally covariant unified field theory, R is charged under the fermion interaction with any other field, including another fermion field.

Therefore eq. (11) is :

$$\hat{H} \psi = \hat{H} \psi \quad (20)$$

$$\hat{H} = \frac{\hbar^2 R}{m} + V + \frac{mc^2}{2\hbar^2 R} \hat{p}^2 \quad (21)$$

$$- \frac{mc^2}{4\hbar^4 R} \left(\hat{\sigma} \cdot \hat{p} \left(\frac{(\hat{E} - V)}{R} \hat{\sigma} \cdot \hat{p} \psi \right) \right)$$

$$\psi = \psi^R_s \quad (22)$$