

## 75(12): The H Atom

The orbitals are well known to be:

$$\psi(r, \theta, \phi) = R_{nl}(r) Y(\theta, \phi) \quad - (1)$$

where  $Y(\theta, \phi)$  are spherical harmonics and

$$R_{nl}(r) = - \left( \frac{2}{na} \right) \left[ \frac{(n-l-1)!}{2n[(n+l)!]^3} \right] \rho^l L_{n+l}^{2l+1}(\rho) e^{-\rho/2} \quad - (2)$$

and  $\rho = \frac{2r}{na}$ ,  $- (3)$

where  $a$  is the Bohr radius. For example, with the fact that  $L_{n+l}^{2l+1}(\rho)$  are the associated Laguerre polynomials, the following is a table of the first few orbitals' radial part:

$n$	$l$	$R_{nl}(r)$
1	0 (1s)	$2 \left( \frac{1}{a} \right)^{3/2} \exp(-\rho/2)$
2	0 (2s)	$\frac{1}{2\sqrt{2}} \left( \frac{1}{a} \right)^{3/2} (2-\rho) \exp(-\rho/2)$
2	1 (2p)	$\frac{1}{2\sqrt{6}} \left( \frac{1}{a} \right)^{3/2} \rho \exp(-\rho/2)$

The H atom orbital for  $n=1, l=0, m=0$

is therefore:

$$\psi = \frac{1}{(\pi a^3)^{1/2}} \exp\left(-\frac{r}{a}\right) \quad - (3)$$

2) The normalization is:

$$\int \psi^* \psi d\tau = 1 \quad - (4)$$

This is evaluated using the volume element of the spherical polar coordinate:

$$dV = r^2 \sin\theta dr d\theta d\phi. \quad - (5)$$

Therefore:

$$\int \psi^* \psi d\tau = \frac{1}{\pi a^2} \int_0^\infty r^2 e^{-2r/a} \int_0^{2\pi} \int_0^\pi \sin\theta d\theta d\phi dr \quad - (6)$$

where:

$$\int_0^{2\pi} \int_0^\pi \sin\theta d\theta d\phi = 4\pi. \quad - (7)$$

$$\text{So } \int \psi^* \psi d\tau = \frac{4}{a^3} \int_0^\infty r^2 \exp\left(-\frac{2r}{a}\right) dr. \quad - (8)$$

Now use:

$$\int x^2 e^{-dx} dx = e^{-dx} \left( -\frac{x^2}{d} - \frac{2x}{d^2} - \frac{2}{d^3} \right) \quad - (9)$$

to find:

$$\int_0^\infty r^2 e^{-2r/a} dr = \frac{a^3}{4} \quad - (10)$$

Therefore:

$$\int \psi^* \psi d\tau = 1 \quad - (11)$$

Q.E.D.

The gradient of  $\phi$  is:

$$\underline{\nabla} \phi = \frac{\partial \phi}{\partial r} \underline{e}_r + \frac{1}{r} \frac{\partial \phi}{\partial \theta} \underline{e}_\theta + \frac{1}{r \sin \phi} \frac{\partial \phi}{\partial \phi} \underline{e}_\phi \quad (12)$$

so:

$$\hat{p} \phi = -i \hbar \underline{\nabla} \phi \quad (13)$$

Therefore

$$\underline{\nabla} \phi = - \frac{\underline{e}_r}{(\pi a^5)^{1/2}} \exp\left(-\frac{r}{a}\right) \quad (14)$$

and

$$\hat{p} \phi = \frac{i \hbar \underline{e}_r}{(\pi a^5)^{1/2}} \exp\left(-\frac{r}{a}\right) \quad (15)$$

and

$$\frac{\partial \phi}{\partial r} = - \frac{1}{(\pi a^5)^{1/2}} \exp\left(-\frac{r}{a}\right) \quad (16)$$

Now we:

$$\underline{r} = x \underline{i} + y \underline{j} + z \underline{k} \quad (17)$$

where:

$$\left. \begin{aligned} x &= r \sin \theta \cos \phi \\ y &= r \sin \theta \sin \phi \\ z &= r \cos \theta \end{aligned} \right\} \quad (18)$$

$$\left. \begin{aligned} \underline{i} &= \sin \theta \cos \phi \underline{e}_r + \cos \theta \cos \phi \underline{e}_\theta - \sin \phi \underline{e}_\phi \\ \underline{j} &= \sin \theta \sin \phi \underline{e}_r + \cos \theta \sin \phi \underline{e}_\theta + \cos \phi \underline{e}_\phi \\ \underline{k} &= \cos \theta \underline{e}_r - \sin \theta \underline{e}_\theta \end{aligned} \right\} \quad (19)$$

To find that:

$$\begin{aligned} \underline{r} &= r (\cos^2 \theta + \sin^2 \theta \cos^2 \phi + \sin^2 \theta \sin^2 \phi) \underline{e}_r \\ &\quad + r (-\cos \theta \sin \theta + \sin \theta \cos \theta \cos^2 \phi + \sin \theta \cos \theta \sin^2 \phi) \underline{e}_\theta \\ &\quad + r (-\cos \theta \sin \theta + \sin \theta \cos \theta \cos^2 \phi + \sin \theta \cos \theta \sin^2 \phi) \underline{e}_\phi \end{aligned}$$

$$+ r(-\sin\theta \cos\phi \sin\phi + \sin\theta \sin\phi \cos\phi) \underline{e}_\phi \\ = r \underline{e}_r \quad - (20)$$

So:

$$\hat{p}_r \psi = \frac{i\hbar}{(\pi a^3)^{1/2}} \exp\left(-\frac{r}{a}\right) \underline{e}_r \quad - (21)$$

$$\underline{r} = r \underline{e}_r \quad - (22)$$

Therefore it is convenient to work out:

$$\begin{aligned} [r, \frac{\partial}{\partial r}] \psi &= r \frac{\partial \psi}{\partial r} - \frac{\partial}{\partial r} (r \psi) \\ &= r \frac{\partial \psi}{\partial r} - r \frac{\partial \psi}{\partial r} - \psi \quad - (23) \end{aligned}$$

and  $\{r, \frac{\partial}{\partial r}\} \psi = 2r \frac{\partial \psi}{\partial r} + \psi \quad - (24)$

For this purely radial wave function:

$$[r, \hat{p}_r] \psi = i\hbar \psi \quad - (25)$$

and  $\{r, \hat{p}_r\} \psi = -i\hbar \left( \psi + 2r \frac{\partial \psi}{\partial r} \right) \quad - (26)$

if  $[r^2, \hat{p}_r^2] \psi = 2i\hbar \{r, \hat{p}_r\} \psi \quad - (27)$

$$= 2\hbar^2 \left( \psi + 2r \frac{\partial \psi}{\partial r} \right)$$

the expectation value can now be worked out.