

172(8): The Wave Form of the EEE Fermion Equation

The EEE fermion equation is:

$$\sigma^0 \hat{E} \psi - c \sigma^3 (\hat{p}_x \psi \sigma^1 - \hat{p}_y \psi \sigma^2 + \hat{p}_z \psi \sigma^3) = mc^2 \sigma^1 \psi \quad (1)$$

where:

$$\psi = \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix}, \quad \sigma^0 = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad \sigma^1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad (2)$$

$$\sigma^2 = \begin{bmatrix} 0 & i \\ -i & 0 \end{bmatrix}, \quad \sigma^3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}.$$

Written out in full, eq. (1) is:

$$\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \hat{E} \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} - c \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \left(\hat{p}_x \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} - \hat{p}_y \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} \begin{bmatrix} 0 & i \\ -i & 0 \end{bmatrix} \right. \\ \left. + \hat{p}_z \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \right) = mc^2 \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} \quad (3)$$

i.e.

$$\hat{E} \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} - c \hat{p}_x \begin{bmatrix} \psi_2^R & \psi_1^R \\ -\psi_2^L & -\psi_1^L \end{bmatrix} + c \hat{p}_y \begin{bmatrix} -i\psi_2^R & i\psi_1^R \\ i\psi_2^L & -i\psi_1^L \end{bmatrix} - c \hat{p}_z \begin{bmatrix} \psi_1^R & -\psi_2^R \\ -\psi_1^L & \psi_2^L \end{bmatrix} = mc^2 \begin{bmatrix} \psi_1^L & \psi_2^L \\ \psi_1^R & \psi_2^R \end{bmatrix} \quad (4)$$

Eq. (4) is equivalent to:

$$\left. \begin{aligned} \hat{E} \psi_1^R - c (\hat{p}_z \psi_1^R + (\hat{p}_x + i\hat{p}_y) \psi_2^R) &= mc^2 \psi_1^L \\ \hat{E} \psi_2^R - c ((\hat{p}_x - i\hat{p}_y) \psi_1^R - \hat{p}_z \psi_2^R) &= mc^2 \psi_2^L \\ \hat{E} \psi_1^L + c (\hat{p}_z \psi_1^L + (\hat{p}_x + i\hat{p}_y) \psi_2^L) &= mc^2 \psi_1^R \\ \hat{E} \psi_2^L + c ((\hat{p}_x - i\hat{p}_y) \psi_1^L - \hat{p}_z \psi_2^L) &= mc^2 \psi_2^R \end{aligned} \right\} \quad (5)$$

2) Eq. (5) is equivalent to the two equations:

$$\left(\hat{E} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} - c \begin{bmatrix} \hat{p}_z & \hat{p}_x + i\hat{p}_y \\ \hat{p}_x - i\hat{p}_y & -\hat{p}_z \end{bmatrix} \right) \begin{bmatrix} \psi_1^R \\ \psi_2^R \end{bmatrix} = mc^2 \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \psi_1^L \\ \psi_2^L \end{bmatrix} \quad (6)$$

and

$$\left(\hat{E} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} + c \begin{bmatrix} \hat{p}_z & \hat{p}_x + i\hat{p}_y \\ \hat{p}_x - i\hat{p}_y & -\hat{p}_z \end{bmatrix} \right) \begin{bmatrix} \psi_1^L \\ \psi_2^L \end{bmatrix} = mc^2 \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \psi_1^R \\ \psi_2^R \end{bmatrix} \quad (7)$$

which may be written in condensed format as:

$$(\hat{E} - c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^R = mc^2 \phi^L \quad (8)$$

$$(\hat{E} + c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^L = mc^2 \phi^R \quad (9)$$

From eqs. (8) and (9):

$$(\hat{E} - c \underline{\sigma} \cdot \underline{\hat{p}})(\hat{E} + c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^L = m^2 c^4 \phi^L \quad (10)$$

$$(\hat{E} - c \underline{\sigma} \cdot \underline{\hat{p}})(\hat{E} + c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^R = m^2 c^4 \phi^R \quad (11)$$

$$(\hat{E} + c \underline{\sigma} \cdot \underline{\hat{p}})(\hat{E} - c \underline{\sigma} \cdot \underline{\hat{p}}) \phi^R = m^2 c^4 \phi^R \quad (12)$$

Now use: $\hat{E} = -i\hbar \frac{\partial}{\partial t}$, $\hat{p} = -i\hbar \underline{\nabla}$ (12)

$$\square = \partial^\mu \partial_\mu = \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2 \quad (13)$$

to find that:

$$\left. \begin{aligned} (\square + (mc/\hbar)^2) \psi_1^R &= 0 \\ (\square + (mc/\hbar)^2) \psi_2^R &= 0 \\ (\square + (mc/\hbar)^2) \psi_1^L &= 0 \\ (\square + (mc/\hbar)^2) \psi_2^L &= 0 \end{aligned} \right\} \quad (14)$$

$$\left(\square + \left(\frac{mc}{\hbar} \right)^2 \right) \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} = 0 \quad - (15)$$

which is the wave form of the ECE fermion equation.
 eq. (1) is a factorization of eq. (15) using the same eigenfunction. Eq. (15) is a limit of the ECE wave equation.

$$(\square + R) \psi_\mu^a = 0 \quad - (16)$$

with

$$\psi_\mu^a = \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix}, \quad R = \left(\frac{mc}{\hbar} \right)^2 \quad - (17)$$

Eq. (16) is equivalent to the tetrad postulate of Cartan:

$$D_\mu \psi_\mu^a = 0 \quad - (18)$$

The above formalism is based on a Cartan geometry and is valid in the general spacetime in $u(2)$ representation space.

It is simpler and more powerful than the chiral presentation of the Dirac equation. This is:

$$(\gamma^\mu \hat{p}_\mu - mc) \psi_D = 0 \quad - (19)$$

here:

$$4) \quad \gamma^\mu \hat{p}_\mu = \gamma^0 \hat{p}_0 + \gamma^i \hat{p}_i \quad - (20)$$

Here:

$$\gamma^0 = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix}, \quad \gamma^1 = \begin{bmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix}, \quad \gamma^2 = \begin{bmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ -i & 0 & 0 & 0 \end{bmatrix}, \quad \gamma^3 = \begin{bmatrix} 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{bmatrix} \quad - (21)$$

In condensed notation, eq. (21) is:

$$\gamma^0 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad \gamma^i = \begin{bmatrix} 0 & -\sigma^i \\ \sigma^i & 0 \end{bmatrix} \quad - (22)$$

So eq. (19) is:

$$\begin{aligned} & \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \begin{bmatrix} \psi_1^R \\ \psi_2^R \\ \psi_1^L \\ \psi_2^L \end{bmatrix} - c \hat{p}_x \begin{bmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \psi_1^R \\ \psi_2^R \\ \psi_1^L \\ \psi_2^L \end{bmatrix} - c \hat{p}_y \begin{bmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ -i & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \psi_1^R \\ \psi_2^R \\ \psi_1^L \\ \psi_2^L \end{bmatrix} \\ & - c \hat{p}_z \begin{bmatrix} 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{bmatrix} \begin{bmatrix} \psi_1^R \\ \psi_2^R \\ \psi_1^L \\ \psi_2^L \end{bmatrix} = mc^2 \begin{bmatrix} \psi_1^R \\ \psi_2^R \\ \psi_1^L \\ \psi_2^L \end{bmatrix} \quad - (23) \end{aligned}$$

where the Dirac spinor is:

$$\psi_0 = \begin{bmatrix} \psi_1^R \\ \psi_2^R \\ \psi_1^L \\ \psi_2^L \end{bmatrix} \quad - (24)$$

The wave format of eq. (19) is:

$$\left(\square + \left(mc / \hbar \right)^2 \right) \psi_0 = 0 \quad - (25)$$

and is equivalent to eq. (15).

5) The ECE Spinor is :

$$\psi = \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} \quad - (26)$$

and is a Catantetrad in Su(2) rep space.

Being a tetrad, it can be written in the general spacetime with torsion and curvature. On the other hand the Dirac spinor is not a tetrad and cannot be written in the general spacetime. The ECE spinor is part of a generally covariant unified field theory, but the Dirac spinor is not. The way that the Dirac equation (19) is usually written in the general spacetime is to attempt to generalize $\gamma^\mu p_\mu$, leaving ψ_0 and m unchanged.

Eq. (1) can be immediately generalized to spacetime with torsion and curvature by writing it

$$\hat{E} \psi - c \sigma^3 (\hat{p}_x \psi \sigma^1 - \hat{p}_y \psi \sigma^2 + \hat{p}_z \psi \sigma^3) = \pm R^{1/2} \sigma^1 \psi \quad - (27)$$